

Some mathematical insights on DMFT applied to the Hubbard model

Well-posedness (pedagogically)

S. Perrin-Roussel^{1,2}, É. Cancès^{1,2}
Alfred Kirsch^{1,2}

¹CERMICS, École des Ponts

²MATERIALS, INRIA

MSQM workshop, January 12

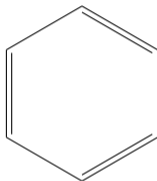
SIMONS
FOUNDATION



Inria



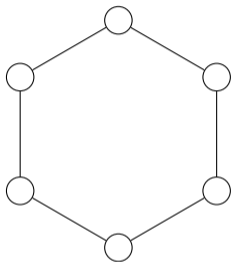
- 1 The Hubbard and Anderson impurity models
- 2 Green's functions and derivatives
- 3 Dynamical Mean-Field Theory hand-wavy
- 4 Well-posedness of IPT-DMFT



Interacting model for the π -electrons

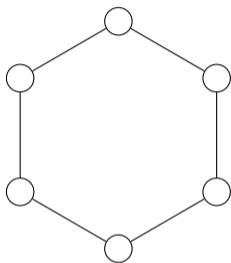
The finite Hubbard model in one slide

- Start with a *graph*
 $\mathcal{G} = (\Lambda, E)$.



(a) The Pariser-Parr-Pople model of benzene : the Hubbard model on the C_6 graph.

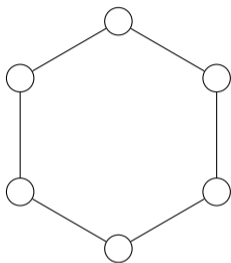
The finite Hubbard model in one slide



(a) The Pariser-Parr-Pople model of benzene : the Hubbard model on the C_6 graph.

- Start with a *graph*
 $\mathcal{G} = (\Lambda, E)$.
- One-site Fock space
 $\mathcal{F}_1 = \text{Vect}(|0\rangle, |\uparrow\rangle, |\downarrow\rangle, |\uparrow\downarrow\rangle)$

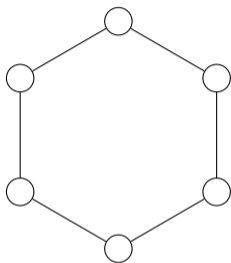
The finite Hubbard model in one slide



(a) The Pariser-Parr-Pople model of benzene : the Hubbard model on the C_6 graph.

- Start with a *graph*
 $\mathcal{G} = (\Lambda, E)$.
- One-site Fock space
 $\mathcal{F}_1 = \text{Vect}(|0\rangle, |\uparrow\rangle, |\downarrow\rangle, |\uparrow\downarrow\rangle)$
- *Sites are distinguishable* :
 $\mathcal{F}_H = \bigotimes_{i \in \Lambda} \mathcal{F}_i$

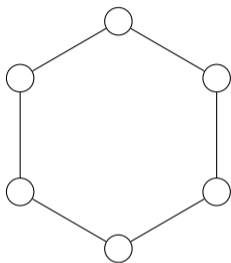
The finite Hubbard model in one slide



(a) The Pariser-Parr-Pople model of benzene : the Hubbard model on the C_6 graph.

- Start with a *graph*
 $\mathcal{G} = (\Lambda, E)$.
- One-site Fock space
 $\mathcal{F}_1 = \text{Vect}(|0\rangle, |\uparrow\rangle, |\downarrow\rangle, |\uparrow\downarrow\rangle)$
- Sites are *distinguishable* :
 $\mathcal{F}_H = \bigotimes_{i \in \Lambda} \mathcal{F}_i$
- Electrons jump :
 $\hat{H}^0 = \sum_{i,j \in E} h_{i,j} \hat{a}_{i,\sigma}^\dagger \hat{a}_{j,\sigma'}$

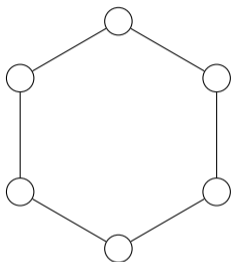
The finite Hubbard model in one slide



(a) The Pariser-Parr-Pople model of benzene : the Hubbard model on the C_6 graph.

- Start with a *graph*
 $\mathcal{G} = (\Lambda, E)$.
- One-site Fock space
 $\mathcal{F}_1 = \text{Vect}(|0\rangle, |\uparrow\rangle, |\downarrow\rangle, |\uparrow\downarrow\rangle)$
- Sites are *distinguishable* :
 $\mathcal{F}_H = \bigotimes_{i \in \Lambda} \mathcal{F}_i$
- Electrons jump :
 $\hat{H}^0 = \sum_{i,j \in E} h_{i,j} \hat{a}_{i,\sigma}^\dagger \hat{a}_{j,\sigma'}$
- Electrons repel *locally*
 $\hat{H}^1 = \sum_{i \in \Lambda} U_i \hat{n}_{i,\uparrow} \hat{n}_{i,\downarrow}$

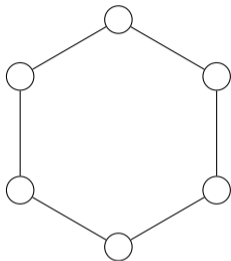
The finite Hubbard model in one slide



(a) The Pariser-Parr-Pople model of benzene : the Hubbard model on the C_6 graph.

- Start with a *graph*
 $\mathcal{G} = (\Lambda, E)$.
- One-site Fock space
 $\mathcal{F}_1 = \text{Vect}(|0\rangle, |\uparrow\rangle, |\downarrow\rangle, |\uparrow\downarrow\rangle)$
- Sites are *distinguishable* :
 $\mathcal{F}_H = \bigotimes_{i \in \Lambda} \mathcal{F}_i$
- Electrons jump :
 $\hat{H}^0 = \sum_{i,j \in E} h_{i,j} \hat{a}_{i,\sigma}^\dagger \hat{a}_{j,\sigma'}$
- Electrons repel *locally*
 $\hat{H}^1 = \sum_{i \in \Lambda} U_i \hat{n}_{i,\uparrow} \hat{n}_{i,\downarrow}$
- The Hubbard Hamiltonian
 $\hat{H}_H = \hat{H}^0 + \hat{H}^1$

The finite Hubbard model in one slide

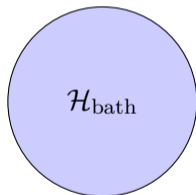
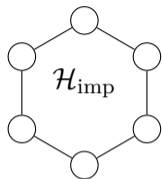


(a) The Pariser-Parr-Pople model of benzene : the Hubbard model on the C_6 graph.

Analytic solutions : [Lieb, 2001]

- Start with a *graph*
 $\mathcal{G} = (\Lambda, E)$.
- One-site Fock space
 $\mathcal{F}_1 = \text{Vect}(|0\rangle, |\uparrow\rangle, |\downarrow\rangle, |\uparrow\downarrow\rangle)$
- Sites are *distinguishable* :
 $\mathcal{F}_H = \bigotimes_{i \in \Lambda} \mathcal{F}_i$
- Electrons jump :
 $\hat{H}^0 = \sum_{i,j \in E} h_{i,j} \hat{a}_{i,\sigma}^\dagger \hat{a}_{j,\sigma'}$
- Electrons repel *locally*
 $\hat{H}^1 = \sum_{i \in \Lambda} U_i \hat{n}_{i,\uparrow} \hat{n}_{i,\downarrow}$
- The Hubbard Hamiltonian
 $\hat{H}_H = \hat{H}^0 + \hat{H}^1$

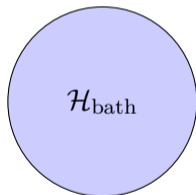
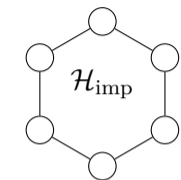
The finite Anderson impurity model in one slide (AIM)



- Bath Fock space:
 $\mathcal{F}_{\text{bath}} = \mathcal{F}(\mathcal{H}_{\text{bath}})$

(a) The Anderson impurity model : an impurity (Hubbard like) and an electronic bath (conducting electrons, σ -electrons)

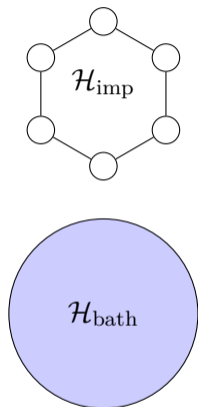
The finite Anderson impurity model in one slide (AIM)



- Bath Fock space:
 $\mathcal{F}_{\text{bath}} = \mathcal{F}(\mathcal{H}_{\text{bath}})$
- AIM Fock space:
 $\mathcal{F}_{\text{AIM}} = \mathcal{F}_{\text{imp}} \otimes \mathcal{F}_{\text{bath}},$
 $\mathcal{F}_{\text{imp}} = \mathcal{F}_H$

(a) The Anderson impurity model : an impurity (Hubbard like) and an electronic bath (conducting electrons, σ -electrons)

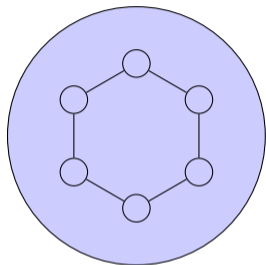
The finite Anderson impurity model in one slide (AIM)



(a) The Anderson impurity model : an impurity (Hubbard like) and an electronic bath (conducting electrons, σ -electrons)

- Bath Fock space:
 $\mathcal{F}_{\text{bath}} = \mathcal{F}(\mathcal{H}_{\text{bath}})$
- AIM Fock space:
 $\mathcal{F}_{\text{AIM}} = \mathcal{F}_{\text{imp}} \otimes \mathcal{F}_{\text{bath}},$
 $\mathcal{F}_{\text{imp}} = \mathcal{F}_H$
- Bath has energy levels:
 $\hat{H}_{\text{bath}}^0 = \sum_{k \in \text{bath}} \epsilon_k \hat{n}_k$

The finite Anderson impurity model in one slide (AIM)

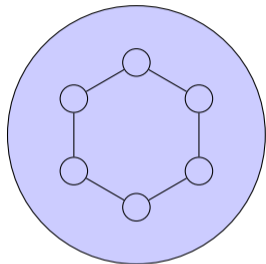


$$\mathcal{H}_{\text{AIM}} = \mathcal{H}_{\text{imp}} \oplus \mathcal{H}_{\text{bath}}$$

(a) The Anderson impurity model : an impurity (Hubbard like) and an electronic bath (conducting electrons, σ -electrons)

- Bath Fock space:
 $\mathcal{F}_{\text{bath}} = \mathcal{F}(\mathcal{H}_{\text{bath}})$
- AIM Fock space:
 $\mathcal{F}_{\text{AIM}} = \mathcal{F}_{\text{imp}} \otimes \mathcal{F}_{\text{bath}},$
 $\mathcal{F}_{\text{imp}} = \mathcal{F}_H$
- Bath has energy levels:
 $\hat{H}_{\text{bath}}^0 = \sum_{k \in \text{bath}} \epsilon_k \hat{n}_k$
- Bath interacts with impurity:
 $\hat{H}_{\text{int}} = \sum_{k \in \text{bath}}^{i \in \Lambda} V_{k,i} \hat{a}_k^\dagger \hat{a}_{i,\sigma}$

The finite Anderson impurity model in one slide (AIM)

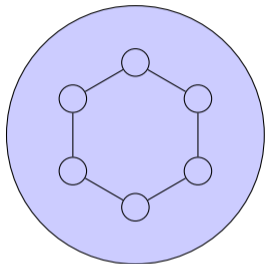


$$\mathcal{H}_{\text{AIM}} = \mathcal{H}_{\text{imp}} \oplus \mathcal{H}_{\text{bath}}$$

(a) The Anderson impurity model : an impurity (Hubbard like) and an electronic bath (conducting electrons, σ -electrons)

- Bath Fock space:
 $\mathcal{F}_{\text{bath}} = \mathcal{F}(\mathcal{H}_{\text{bath}})$
- AIM Fock space:
 $\mathcal{F}_{\text{AIM}} = \mathcal{F}_{\text{imp}} \otimes \mathcal{F}_{\text{bath}},$
 $\mathcal{F}_{\text{imp}} = \mathcal{F}_H$
- Bath has energy levels:
 $\hat{H}_{\text{bath}}^0 = \sum_{k \in \text{bath}} \epsilon_k \hat{n}_k$
- Bath interacts with impurity:
 $\hat{H}_{\text{int}} = \sum_{k \in \text{bath}}^{i \in \Lambda} V_{k,i} \hat{a}_k^\dagger \hat{a}_{i,\sigma}$
- AIM non interacting:
 $\hat{H}_{\text{AIM}}^0 = \hat{H}_H^0 + \hat{H}_{\text{bath}}^0 + \hat{H}_{\text{int}}$

The finite Anderson impurity model in one slide (AIM)



$$\mathcal{H}_{\text{AIM}} = \mathcal{H}_{\text{imp}} \oplus \mathcal{H}_{\text{bath}}$$

(a) The Anderson impurity model : an impurity (Hubbard like) and an electronic bath (conducting electrons, σ -electrons)

- Bath Fock space:
 $\mathcal{F}_{\text{bath}} = \mathcal{F}(\mathcal{H}_{\text{bath}})$
- AIM Fock space:
 $\mathcal{F}_{\text{AIM}} = \mathcal{F}_{\text{imp}} \otimes \mathcal{F}_{\text{bath}},$
 $\mathcal{F}_{\text{imp}} = \mathcal{F}_H$
- Bath has energy levels:
 $\hat{H}_{\text{bath}}^0 = \sum_{k \in \text{bath}} \epsilon_k \hat{n}_k$
- Bath interacts with impurity:
 $\hat{H}_{\text{int}} = \sum_{k \in \text{bath}}^{i \in \Lambda} V_{k,i} \hat{a}_k^\dagger \hat{a}_{i,\sigma}$
- AIM non interacting:
 $\hat{H}_{\text{AIM}}^0 = \hat{H}_H^0 + \hat{H}_{\text{bath}}^0 + \hat{H}_{\text{int}}$
- AIM Hamiltonian:
 $\hat{H}_{\text{AIM}} = \hat{H}_{\text{AIM}}^0 + \hat{H}_H^1$

Green's functions are not
Green's functions.

Green's functions are not
always *Green's functions*.

Quantum Green's functions are not
always *mathematical Green's functions*.

Quantum physics Green's functions are not
always *mathematical Green's functions*.

Quantum physics Green's functions (Propagators) are not always *mathematical Green's functions* (Fundamental solution associated to a Linear Differential Operator).

Heisenberg picture : $\mathbb{H}(\mathcal{O})(t) = e^{-it\hat{H}}\mathcal{O}e^{-it\hat{H}}$.

State is a linear form on operators $\Gamma(\mathcal{O}) = \text{Tr}(\rho\mathcal{O})$, $\rho^2 \leq \rho$, equilibrium $[\rho, \hat{H}] = 0$

Definition (Green's functions)

The *one-body time-ordered Green's function* \tilde{G} of a quantum system \mathcal{H} , \hat{H} in a state Γ is the *matrix-valued* function with

$$i\tilde{G}_{i\sigma,j\sigma'}(t) = \Theta(t)\Gamma(\mathbb{H}(\hat{a}_{i,\sigma})(t)\hat{a}_{j,\sigma'}^\dagger) + \Theta(-t)\Gamma(\mathbb{H}(\hat{a}_{j,\sigma'}^\dagger)(t)\hat{a}_{i,\sigma})$$

Quantum physics Green's functions

Heisenberg picture : $\mathbb{H}(\mathcal{O})(t) = e^{-it\hat{H}}\mathcal{O}e^{-it\hat{H}}$.

State is a linear form on operators $\Gamma(\mathcal{O}) = \text{Tr}(\rho\mathcal{O})$, $\rho^2 \leq \rho$, equilibrium $[\rho, \hat{H}] = 0$

Definition (Green's functions)

The *one-body time-ordered Green's function* \tilde{G} of a quantum system \mathcal{H} , \hat{H} in a state Γ is the *matrix-valued function* with

$$i\tilde{G}_{i\sigma,j\sigma'}(t) = \underbrace{\Theta(t)\Gamma(\mathbb{H}(\hat{a}_{i,\sigma})(t)\hat{a}_{j,\sigma'}^\dagger)}_{\text{particle}} + \Theta(-t)\Gamma(\mathbb{H}(\hat{a}_{j,\sigma'}^\dagger)(t)\hat{a}_{i,\sigma})$$

Quantum physics Green's functions

Heisenberg picture : $\mathbb{H}(\mathcal{O})(t) = e^{-it\hat{H}}\mathcal{O}e^{-it\hat{H}}$.

State is a linear form on operators $\Gamma(\mathcal{O}) = \text{Tr}(\rho\mathcal{O})$, $\rho^2 \leq \rho$, equilibrium $[\rho, \hat{H}] = 0$

Definition (Green's functions)

The *one-body time-ordered Green's function* \tilde{G} of a quantum system \mathcal{H}, \hat{H} in a state Γ is the *matrix-valued function* with

$$i\tilde{G}_{i\sigma,j\sigma'}(t) = \underbrace{\Theta(t)\Gamma(\mathbb{H}(\hat{a}_{i,\sigma})(t)\hat{a}_{j,\sigma'}^\dagger)}_{\text{particle}} + \underbrace{\Theta(-t)\Gamma(\mathbb{H}(\hat{a}_{j,\sigma'}^\dagger)(t)\hat{a}_{i,\sigma})}_{\text{hole}}$$

Quantum physics Green's functions

Heisenberg picture : $\mathbb{H}(\mathcal{O})(t) = e^{-it\hat{H}}\mathcal{O}e^{-it\hat{H}}$.

State is a linear form on operators $\Gamma(\mathcal{O}) = \text{Tr}(\rho\mathcal{O})$, $\rho^2 \leq \rho$, equilibrium $[\rho, \hat{H}] = 0$

Definition (Green's functions)

The *one-body time-ordered Green's function* \tilde{G} of a quantum system \mathcal{H} , \hat{H} in a state Γ is the *matrix-valued function* with

$$i\tilde{G}_{i\sigma,j\sigma'}(t) = \underbrace{\Theta(t)\Gamma(\mathbb{H}(\hat{a}_{i,\sigma})(t)\hat{a}_{j,\sigma'}^\dagger)}_{\text{particle}} + \underbrace{\Theta(-t)\Gamma(\mathbb{H}(\hat{a}_{j,\sigma'}^\dagger)(t)\hat{a}_{i,\sigma})}_{\text{hole}}$$

- Quantum physics Green's functions are *explicitly* defined.

Quantum physics Green's functions

Heisenberg picture : $\mathbb{H}(\mathcal{O})(t) = e^{-it\hat{H}}\mathcal{O}e^{-it\hat{H}}$.

State is a linear form on operators $\Gamma(\mathcal{O}) = \text{Tr}(\rho\mathcal{O})$, $\rho^2 \leq \rho$, equilibrium $[\rho, \hat{H}] = 0$

Definition (Green's functions)

The *one-body time-ordered Green's function* \tilde{G} of a quantum system \mathcal{H} , \hat{H} in a state Γ is the *matrix-valued* function with

$$i\tilde{G}_{i\sigma,j\sigma'}(t) = \underbrace{\Theta(t)\Gamma(\mathbb{H}(\hat{a}_{i,\sigma})(t)\hat{a}_{j,\sigma'}^\dagger)}_{\text{particle}} + \underbrace{\Theta(-t)\Gamma(\mathbb{H}(\hat{a}_{j,\sigma'}^\dagger)(t)\hat{a}_{i,\sigma})}_{\text{hole}}$$

- Quantum physics Green's functions are *explicitly* defined.
- *Enough* to compute many observables : average energy (Galitski-Migdal), conduction behaviour, Chern number, etc.

Quantum physics Green's functions

Heisenberg picture : $\mathbb{H}(\mathcal{O})(t) = e^{-it\hat{H}}\mathcal{O}e^{-it\hat{H}}$.

State is a linear form on operators $\Gamma(\mathcal{O}) = \text{Tr}(\rho\mathcal{O})$, $\rho^2 \leq \rho$, equilibrium $[\rho, \hat{H}] = 0$

Definition (Green's functions)

The *one-body time-ordered Green's function* \tilde{G} of a quantum system \mathcal{H}, \hat{H} in a state Γ is the *matrix-valued* function with

$$i\tilde{G}_{i\sigma, j\sigma'}(t) = \underbrace{\Theta(t)\Gamma(\mathbb{H}(\hat{a}_{i,\sigma})(t)\hat{a}_{j,\sigma'}^\dagger)}_{\text{particle}} + \underbrace{\Theta(-t)\Gamma(\mathbb{H}(\hat{a}_{j,\sigma'}^\dagger)(t)\hat{a}_{i,\sigma})}_{\text{hole}}$$

- Quantum physics Green's functions are *explicitly* defined.
- *Enough* to compute many observables : average energy (Galitski-Migdal), conduction behaviour, Chern number, etc.
- *Experimentally* "measurable" : ARPES (see Lucia's talk)

If $\dim(\mathcal{H})$ is finite, \tilde{G} is oscillatory (Lehmann-Källén's representation).

If $\dim(\mathcal{H})$ is finite, \tilde{G} is oscillatory (Lehmann-Källén's representation).

Definition (Generalized Fourier transform)

The Generalized Fourier transform [Titchmarsh,1948] of \tilde{G} is the map G defined for $z \in \mathbb{C}_+$ by

$$G(z) = \int_{\mathbb{R}^+} e^{izt} \tilde{G}(t) dt + \int_{\mathbb{R}^-} e^{i\bar{z}t} \tilde{G}(t) dt$$

Well-defined and invertible.

– $G : \mathbb{C}_+ \rightarrow \mathbb{C}_+$ and *analytic* :

If $\dim(\mathcal{H})$ is finite, \tilde{G} is oscillatory (Lehmann-Källén's representation).

Definition (Generalized Fourier transform)

The Generalized Fourier transform [Titchmarsh,1948] of \tilde{G} is the map G defined for $z \in \mathbb{C}_+$ by

$$G(z) = \int_{\mathbb{R}^+} e^{izt} \tilde{G}(t) dt + \int_{\mathbb{R}^-} e^{i\bar{z}t} \tilde{G}(t) dt$$

Well-defined and invertible.

$-G : \mathbb{C}_+ \rightarrow \mathbb{C}_+$ and *analytic* : Herglotz functions.

a.k.a

If $\dim(\mathcal{H})$ is finite, \tilde{G} is oscillatory (Lehmann-Källén's representation).

Definition (Generalized Fourier transform)

The Generalized Fourier transform [Titchmarsh,1948] of \tilde{G} is the map G defined for $z \in \mathbb{C}_+$ by

$$G(z) = \int_{\mathbb{R}^+} e^{izt} \tilde{G}(t) dt + \int_{\mathbb{R}^-} e^{i\bar{z}t} \tilde{G}(t) dt$$

Well-defined and invertible.

$-G : \mathbb{C}_+ \rightarrow \mathbb{C}_+$ and *analytic* : Herglotz functions.

a.k.a Pick,

If $\dim(\mathcal{H})$ is finite, \tilde{G} is oscillatory (Lehmann-Källén's representation).

Definition (Generalized Fourier transform)

The Generalized Fourier transform [Titchmarsh,1948] of \tilde{G} is the map G defined for $z \in \mathbb{C}_+$ by

$$G(z) = \int_{\mathbb{R}^+} e^{izt} \tilde{G}(t) dt + \int_{\mathbb{R}^-} e^{i\bar{z}t} \tilde{G}(t) dt$$

Well-defined and invertible.

– $G : \mathbb{C}_+ \rightarrow \mathbb{C}_+$ and *analytic* : Herglotz functions.

a.k.a Pick, Nevanlinna,

If $\dim(\mathcal{H})$ is finite, \tilde{G} is oscillatory (Lehmann-Källén's representation).

Definition (Generalized Fourier transform)

The Generalized Fourier transform [Titchmarsh,1948] of \tilde{G} is the map G defined for $z \in \mathbb{C}_+$ by

$$G(z) = \int_{\mathbb{R}^+} e^{izt} \tilde{G}(t) dt + \int_{\mathbb{R}^-} e^{i\bar{z}t} \tilde{G}(t) dt$$

Well-defined and invertible.

– $G : \mathbb{C}_+ \rightarrow \mathbb{C}_+$ and *analytic* : Herglotz functions.

a.k.a Pick, Nevanlinna, Riesz,

If $\dim(\mathcal{H})$ is finite, \tilde{G} is oscillatory (Lehmann-Källén's representation).

Definition (Generalized Fourier transform)

The Generalized Fourier transform [Titchmarsh,1948] of \tilde{G} is the map G defined for $z \in \mathbb{C}_+$ by

$$G(z) = \int_{\mathbb{R}^+} e^{izt} \tilde{G}(t) dt + \int_{\mathbb{R}^-} e^{i\bar{z}t} \tilde{G}(t) dt$$

Well-defined and invertible.

– $G : \mathbb{C}_+ \rightarrow \mathbb{C}_+$ and *analytic* : Herglotz functions.

a.k.a Pick, Nevanlinna, Riesz, Weyl,

If $\dim(\mathcal{H})$ is finite, \tilde{G} is oscillatory (Lehmann-Källén's representation).

Definition (Generalized Fourier transform)

The Generalized Fourier transform [Titchmarsh,1948] of \tilde{G} is the map G defined for $z \in \mathbb{C}_+$ by

$$G(z) = \int_{\mathbb{R}^+} e^{izt} \tilde{G}(t) dt + \int_{\mathbb{R}^-} e^{i\bar{z}t} \tilde{G}(t) dt$$

Well-defined and invertible.

– $G : \mathbb{C}_+ \rightarrow \mathbb{C}_+$ and *analytic* : Herglotz functions.

a.k.a Pick, Nevanlinna, Riesz, Weyl, Titchmarsh,

If $\dim(\mathcal{H})$ is finite, \tilde{G} is oscillatory (Lehmann-Källén's representation).

Definition (Generalized Fourier transform)

The Generalized Fourier transform [Titchmarsh,1948] of \tilde{G} is the map G defined for $z \in \mathbb{C}_+$ by

$$G(z) = \int_{\mathbb{R}^+} e^{izt} \tilde{G}(t) dt + \int_{\mathbb{R}^-} e^{i\bar{z}t} \tilde{G}(t) dt$$

Well-defined and invertible.

– $G : \mathbb{C}_+ \rightarrow \mathbb{C}_+$ and *analytic* : Herglotz functions.

a.k.a Pick, Nevanlinna, Riesz, Weyl, Titchmarsh, *R*-function

Non-interacting Green's function

The *non-interacting Green's function* G^0 is the Green's function of $(\mathcal{H}, \hat{H}^0, \Gamma)$.

Non-interacting Green's functions are Green's functions

Assume $\hat{H}^0 = \sum_{i,j} h_{i,j} \hat{a}_i^\dagger \hat{a}_j$. Then, the non-interacting Green's function G^0 is the *resolvent* of h :

$$G^0(z) = (z - h)^{-1}.$$

Non-interacting Green's function

The *non-interacting Green's function* G^0 is the Green's function of $(\mathcal{H}, \hat{H}^0, \Gamma)$.

Non-interacting Green's functions are Green's functions

Assume $\hat{H}^0 = \sum_{i,j} h_{i,j} \hat{a}_i^\dagger \hat{a}_j$. Then, the non-interacting Green's function G^0 is the *resolvent* of h :

$$G^0(z) = (z - h)^{-1}.$$

In the time domain, equivalent to $(i\partial_t - h) \tilde{G}^0 = \delta.$

Definition (Self-energy)

The *self-energy* Σ associated to $(\mathcal{H}, \hat{H}, \Gamma)$ is the map defined for all $z \in \mathbb{C}_+$ by

$$\Sigma(z) = (G^0(z))^{-1} - G(z)^{-1}. \quad (1)$$

- Etymology : (1) $\iff G(z) = (z - (h + \Sigma(z)))^{-1}$

Definition (Self-energy)

The *self-energy* Σ associated to $(\mathcal{H}, \hat{H}, \Gamma)$ is the map defined for all $z \in \mathbb{C}_+$ by

$$\Sigma(z) = (G^0(z))^{-1} - G(z)^{-1}. \quad (1)$$

- Etymology : (1) $\iff G(z) = (z - \underbrace{(h + \Sigma(z))}_{\text{effective } h'})^{-1}$
- $-\Sigma$ is a Herglotz function.

Definition (Self-energy)

The *self-energy* Σ associated to $(\mathcal{H}, \hat{H}, \Gamma)$ is the map defined for all $z \in \mathbb{C}_+$ by

$$\Sigma(z) = (G^0(z))^{-1} - G(z)^{-1}. \quad (1)$$

- Etymology : (1) $\iff G(z) = (z - \underbrace{(h + \Sigma(z))}_{\text{effective } h'})^{-1}$
- $-\Sigma$ is a Herglotz function.
- If $\dim(\mathcal{H})$ is finite (see M.Lindsey's thesis), $\exists a_k \in \mathcal{S}(\mathbb{C})^+, \epsilon_k \in \mathbb{R}$ s.t. $\forall z \in \mathbb{C}_+,$

$$\Sigma(z) = \Sigma_{HF} + \sum_k \frac{1}{z - \epsilon_k} a_k \quad (2)$$

DMFT ustencil : the hybridization function Δ .

DMFT is formulated with Green's functions *blocks* : for an Anderson impurity model we have

$$G_{\text{imp}}^0 = (z - h)_{\text{imp}}^{-1} = \left(z - h_{\text{imp}} - \sum_{k=1}^{\text{Bathsize}} \frac{1}{z - \epsilon_k} V_k V_k^\dagger \right)^{-1} \quad (3)$$

Definition (Hybridization function)

The *hybridization function* Δ associated to an Anderson impurity model is the *matrix-valued* map defined for all $z \in \mathbb{C}_+$ by

$$\Delta(z) = \sum_{k=1}^{\text{Bathsize}} \frac{1}{z - \epsilon_k} V_k V_k^\dagger \quad (4)$$

DMFT ustensil : the hybridization function Δ .

DMFT is formulated with Green's functions *blocks* : for an Anderson impurity model we have

$$G_{\text{imp}}^0 = (z - h)_{\text{imp}}^{-1} = \left(z - h_{\text{imp}} - \sum_{k=1}^{\text{Bathsize}} \frac{1}{z - \epsilon_k} V_k V_k^\dagger \right)^{-1} \quad (3)$$

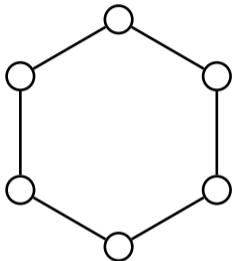
Definition (Hybridization function)

The *hybridization function* Δ associated to an Anderson impurity model is the *matrix-valued* map defined for all $z \in \mathbb{C}_+$ by

$$\Delta(z) = \sum_{k=1}^{\text{Bathsize}} \frac{1}{z - \epsilon_k} V_k V_k^\dagger \quad (4)$$

Δ *fully characterizes* the bath and its coupling to the impurity

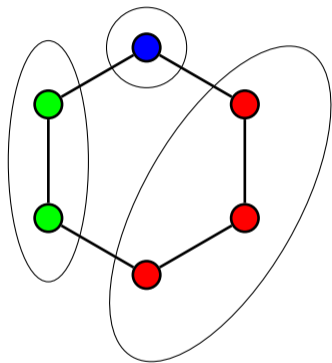
DMFT recipe: from one Hubbard to several small Anderson



$$G = \begin{bmatrix} * & * & * & * & * & * \\ * & * & * & * & * & * \\ * & * & * & * & * & * \\ * & * & * & * & * & * \\ * & * & * & * & * & * \\ * & * & * & * & * & * \end{bmatrix}$$

DMFT recipe : find the Greens function G of a Hubbard model $(\mathcal{F}_H, \hat{H}_H)$ in a state Γ .

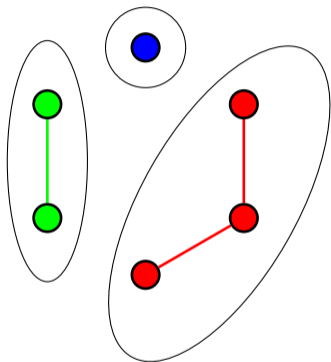
DMFT recipe: from one Hubbard to several small Anderson



$$G = \begin{bmatrix} G_1 & & & & & \\ & * & * & * & & \\ & * & G_2 & * & & \\ & * & * & * & & \\ & & & & G_3 & * \\ & & & & * & * \end{bmatrix}$$

Step 1 : *partition* the vertices $\Lambda = \sqcup_{i=1}^N \Lambda_i$ of the original Hubbard graph $\mathcal{G} = (\Lambda, E)$ and focus on the blocks $(G_i)_{i=1,N}$.

DMFT recipe: from one Hubbard to several small Anderson

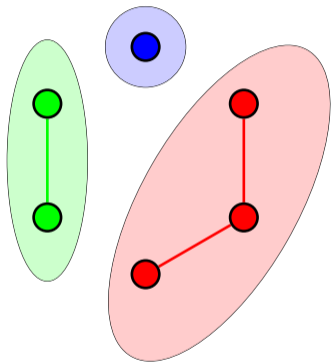


$$G = \begin{bmatrix} G_1 & & & & & \\ & * & * & * & & \\ & * & G_2 & * & & \\ & * & * & * & & \\ & & & & G_3 & * \\ & & & & * & * \end{bmatrix}$$

Step 2 : Define $\mathcal{G}_i = (\Lambda_i, E_i)$, $E_i = \{\{k, l\} \in E, k, l \in \Lambda_i\}$

G of \hat{H}_H with $\mathcal{G}_i \neq G_i$ of \hat{H}_H with \mathcal{G} ! *Even if not interacting !*

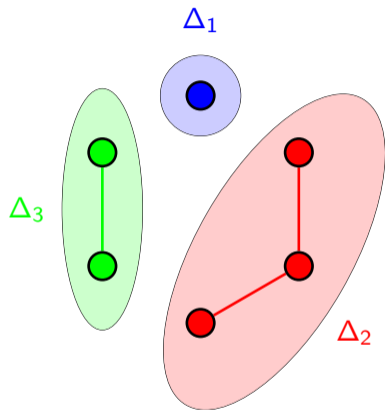
DMFT recipe: from one Hubbard to several small Anderson



$$G = \begin{bmatrix} G_1 & & & & & \\ & * & * & * & & \\ & * & G_2 & * & & \\ & * & * & * & & \\ & & & & G_3 & * \\ & & & & * & * \end{bmatrix}$$

Step 3 : define $G_{\text{imp},i}$ of an Anderson impurity model with an *electronic bath* for each impurity

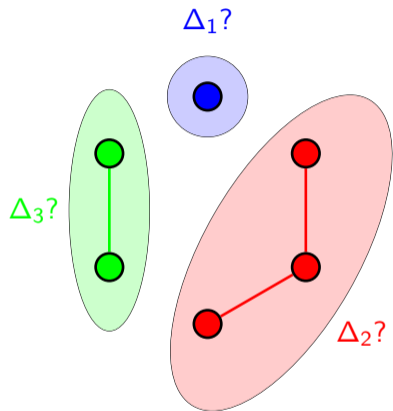
DMFT recipe: from one Hubbard to several small Anderson



$$G = \begin{bmatrix} G_1 & & & & & \\ & * & * & * & & \\ & * & G_2 & * & & \\ & * & * & * & & \\ & & & & G_3 & * \\ & & & & * & * \end{bmatrix}$$

Step 3 : define $G_{\text{imp},i}$ of an Anderson impurity model with an *electronic bath* for each impurity $(\Delta_i)_{i=1,N}$.

DMFT recipe: from one Hubbard to several small Anderson



$$G = \begin{bmatrix} G_1 & & & & & \\ & * & * & * & & \\ & * & G_2 & * & & \\ & * & * & * & & \\ & & & & G_3 & * \\ & & & & * & * \end{bmatrix}$$

Step 3 : define $G_{\text{imp},i}$ of an Anderson impurity model with an *electronic bath* for each impurity $(\Delta_i)_{i=1,N}$. Which one ?

Question : (Δ_i) ?

Self-consistent equation on (Δ_i)

Question : (Δ_i) ?

- First answer : s.t. $G_{\text{imp},i} = G_i$ of \hat{H}_H of \mathcal{G} .

Self-consistent equation on (Δ_i)

Question : (Δ_i) ?

- First answer : s.t. $G_{\text{imp},i} = G_i?$ of \hat{H}_H of \mathcal{G} . **Unknown ! X**

Self-consistent equation on (Δ_i)

Question : (Δ_i) ?

- First answer : s.t. $G_{\text{imp},i} = G_i?$ of \hat{H}_H of \mathcal{G} . **Unknown ! X**
- DFMT answer : s.t. $G_{\text{imp},i} = G_{\text{DMFT},i}$.

Question : (Δ_i) ?

- First answer : s.t. $G_{\text{imp},i} = G_i?$ of \hat{H}_H of \mathcal{G} . **Unknown ! X**
- DFMT answer : s.t. $G_{\text{imp},i} = G_{\text{DMFT},i} \cdot G_{\text{DMFT}}$?
- Requirement : $G_{\text{DMFT}} = G^0$ if not interacting (exact)
 $\iff \Sigma_{\text{DMFT}} = (G^0)^{-1} - G_{\text{DMFT}}^{-1} = 0.$

Question : (Δ_i) ?

- First answer : s.t. $G_{\text{imp},i} = G_i$? of \hat{H}_H of \mathcal{G} . **Unknown ! X**
- DFMT answer : s.t. $G_{\text{imp},i} = G_{\text{DMFT},i}$. G_{DMFT} ?
- Requirement : $G_{\text{DMFT}} = G^0$ if not interacting (exact)
 $\iff \Sigma_{\text{DMFT}} = (G^0)^{-1} - G_{\text{DMFT}}^{-1} = 0$.
- DMFT self-consistency : $G_{\text{DMFT}} = ((G^0)^{-1} - \Sigma_{\text{DMFT}})^{-1}$ with Σ_{DMFT}

$$\Sigma_{\text{DMFT}} = \bigoplus_{i=1}^N \Sigma_{\text{imp},i} \quad (5)$$

The impurity solver : $\Delta_j \rightarrow \Sigma_{\text{imp},j}$

Impurity solver : $\Delta_j \rightarrow \Sigma_j$ for each impurity.

The impurity solver : $\Delta_j \rightarrow \Sigma_{\text{imp},j}$

Impurity solver : $\Delta_j \rightarrow \Sigma_j$ for each impurity.

- One theoretical, many in practical computations : Iterated Perturbation Theory (IPT), continuous-time Monte Carlo, exact diagonalisation etc.

The impurity solver : $\Delta_j \rightarrow \Sigma_{\text{imp},j}$

Impurity solver : $\Delta_j \rightarrow \Sigma_j$ for each impurity.

- One theoretical, many in practical computations : **Iterated Perturbation Theory (IPT)**, continuous-time Monte Carlo, exact diagonalisation etc.

Impurity solver : $\Delta_j \rightarrow \Sigma_j$ for each impurity.

- One theoretical, many in practical computations : **Iterated Perturbation Theory** (IPT), continuous-time Monte Carlo, exact diagonalisation etc.
- Doesn't require a (\mathcal{H}, \hat{H}) formulation (!)

The impurity solver : $\Delta_j \rightarrow \Sigma_{\text{imp},j}$

Impurity solver : $\Delta_j \rightarrow \Sigma_j$ for each impurity.

- One theoretical, many in practical computations : **Iterated Perturbation Theory** (IPT), continuous-time Monte Carlo, exact diagonalisation etc.
- Doesn't require a (\mathcal{H}, \hat{H}) formulation (!)
- Computationally expensive part of DMFT.

The impurity solver : $\Delta_i \rightarrow \Sigma_{\text{imp},i}$

Impurity solver : $\Delta_i \rightarrow \Sigma_i$ for each impurity.

- One theoretical, many in practical computations : **Iterated Perturbation Theory** (IPT), continuous-time Monte Carlo, exact diagonalisation etc.
- Doesn't require a (\mathcal{H}, \hat{H}) formulation (!)
- Computationally expensive part of DMFT.

For each impurity, $\Sigma_i = \text{IPT}(\Delta_i)$

We are looking for Δ_i such that for all $z \in \mathbb{C}_+$,

$$G_{\text{imp},i}(z) = G_{DMFT,i}(z) \tag{6}$$

(7)

(8)

(9)

We are looking for Δ_i such that for all $z \in \mathbb{C}_+$,

$$G_{\text{imp},i}(z) = G_{DMFT,i}(z) \tag{6}$$

$$G_{\text{imp},i}^{-1}(z) = G_{DMFT,i}^{-1}(z) \tag{7}$$

$$\tag{8}$$

$$\tag{9}$$

We are looking for Δ_i such that for all $z \in \mathbb{C}_+$,

$$G_{\text{imp},i}(z) = G_{DMFT,i}(z) \quad (6)$$

$$G_{\text{imp},i}^{-1}(z) = G_{DMFT,i}^{-1}(z) \quad (7)$$

$$\left(\left[(z - h_{\text{AIM},i} - \Sigma_{\text{AIM},i}(z))^{-1} \right]_{\text{imp}} \right)^{-1} = \left(\left[\left(z - h - \bigoplus_{i=1}^N \Sigma_i(z) \right)^{-1} \right]_i \right)^{-1} \quad (8)$$

$$(9)$$

DMFT equations (1)

We are looking for Δ_i such that for all $z \in \mathbb{C}_+$,

$$G_{\text{imp},i}(z) = G_{DMFT,i}(z) \quad (6)$$

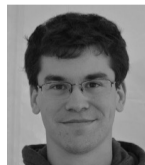
$$G_{\text{imp},i}^{-1}(z) = G_{DMFT,i}^{-1}(z) \quad (7)$$

$$\left(\left[(z - h_{\text{AIM},i} - \Sigma_{\text{AIM},i}(z))^{-1} \right]_{\text{imp}} \right)^{-1} = \left(\left[\left(z - h - \bigoplus_{i=1}^N \Sigma_i(z) \right)^{-1} \right]_i \right)^{-1} \quad (8)$$

(9)



Schur-Levitt
complement !



We are looking for Δ_i such that for all $z \in \mathbb{C}_+$,

$$G_{\text{imp},i}(z) = G_{DMFT,i}(z) \quad (6)$$

$$G_{\text{imp},i}^{-1}(z) = G_{DMFT,i}^{-1}(z) \quad (7)$$

$$\left(\left[(z - h_{\text{AIM},i} - \Sigma_{\text{AIM},i}(z))^{-1} \right]_{\text{imp}} \right)^{-1} = \left(\left[\left(z - h - \bigoplus_{i=1}^N \Sigma_i(z) \right)^{-1} \right]_i \right)^{-1} \quad (8)$$

$$z - h_i - \Sigma_i(z) - \Delta_i(z) = z - h_i - \Sigma_i(z) - h_{i,\bar{i}} \left(z - h_{\bar{i}} - \bigoplus_{j \neq i, j=1}^N \Sigma_j(z) \right)^{-1} h_{i,\bar{i}}^\dagger \quad (9)$$

DMFT equations : for all $i = 1, N$

$$\Delta_i = h_{i,\bar{i}} \left(\cdot - h_{\bar{i}} - \bigoplus_{j \neq i, j=1}^N \Sigma_j \right)^{-1} h_{i,\bar{i}}^\dagger, \quad (\Sigma_i)_{i=1, N} \mapsto \Delta_i \quad (\text{Self-consistent, global})$$

DMFT equations : for all $i = 1, N$

$$\Delta_i = h_{i,\bar{i}} \left(\cdot - h_{\bar{i}} - \bigoplus_{j \neq i, j=1}^N \Sigma_j \right)^{-1} h_{i,\bar{i}}^\dagger, \quad (\Sigma_i)_{i=1, N} \mapsto \Delta_i \quad (\text{Self-consistent, global})$$

$$\Sigma_i = \text{IPT}(\Delta_i), \quad \Delta_i \mapsto \Sigma_i \quad (\text{IPT equation, local})$$

with IPT as an impurity solver.

DMFT equations : for all $i = 1, N$

$$\Delta_i = h_{i,\bar{i}} \left(\cdot - h_{\bar{i}} - \bigoplus_{j \neq i, j=1}^N \Sigma_j \right)^{-1} h_{i,\bar{i}}^\dagger, \quad (\Sigma_i)_{i=1, N} \mapsto \Delta_i \quad (\text{Self-consistent, global})$$

$$\Sigma_i = \text{IPT}(\Delta_i), \quad \Delta_i \mapsto \Sigma_i \quad (\text{IPT equation, local})$$

with IPT as an impurity solver.

DMFT unknowns : $\Delta_i \in \mathfrak{D}_i, \Sigma_i \in \mathfrak{G}_i$

DMFT equations : for all $i = 1, N$

$$\Delta_i = h_{i,\bar{i}} \left(\cdot - h_{\bar{i}} - \bigoplus_{j \neq i, j=1}^N \Sigma_j \right)^{-1} h_{i,\bar{i}}^\dagger, \quad (\Sigma_i)_{i=1, N} \mapsto \Delta_i \quad (\text{Self-consistent, global})$$

$$\Sigma_i = \text{IPT}(\Delta_i), \quad \Delta_i \mapsto \Sigma_i \quad (\text{IPT equation, local})$$

with IPT as an impurity solver.

DMFT unknowns : $\Delta_i \in \mathfrak{D}_i, \Sigma_i \in \mathfrak{G}_i ?$

DMFT equations : for all $i = 1, N$

$$\Delta_i = h_{i,\bar{i}} \left(\cdot - h_{\bar{i}} - \bigoplus_{j \neq i, j=1}^N \Sigma_j \right)^{-1} h_{i,\bar{i}}^\dagger, \quad (\Sigma_i)_{i=1, N} \mapsto \Delta_i \quad (\text{Self-consistent, global})$$

$$\Sigma_i = \text{IPT}(\Delta_i), \quad \Delta_i \mapsto \Sigma_i \quad (\text{IPT equation, local})$$

with IPT as an impurity solver.

DMFT unknowns : $\Delta_i \in \mathfrak{D}_i, \Sigma_i \in \mathfrak{G}_i ?$

Mathematical question : $\mathfrak{D}, \mathfrak{G}$ s. t. well posed and existing solution ?

Self-coherent equation well-posedness for finite bath

Finite bath dimension well-posedness (L. Lin, M. Lindsey, R. Schneider, 2019)

Assume $\forall i = 1, n, \exists C_i \in \mathcal{S}_{N_i}(\mathbb{C}), L_i \in \mathbb{N}, \forall k = 1, L_i, a_k \in \mathcal{S}_{N_i}(\mathbb{C})^+, \epsilon_k \in \mathbb{R}$ s.t.
 $\forall z \in \mathbb{C}^+$,

$$\Sigma_i(z) = C_i + \sum_{k=1}^{L_i} \frac{1}{z - \epsilon_k} a_k \quad (10)$$

Then $\forall i = 1, N, \Delta_i$ is well-defined and there exists $\tilde{L}_i \in \mathbb{N}$,
 $\forall k = 1, \tilde{L}_i, \tilde{a}_k \in \mathcal{S}_{N_i}(\mathbb{C})^+, \tilde{\epsilon}_k \in \mathbb{R}$ s.t. $\forall z \in \mathbb{C}^+$,

$$\Delta_i(z) = \sum_{k=1}^{\tilde{L}_i} \frac{1}{z - \tilde{\epsilon}_k} \tilde{a}_k \quad (11)$$

Self-coherent equation well-posedness for finite bath

Finite bath dimension well-posedness (L. Lin, M. Lindsey, R. Schneider, 2019)

Assume $\forall i = 1, n, \exists C_i \in \mathcal{S}_{N_i}(\mathbb{C}), L_i \in \mathbb{N}, \forall k = 1, L_i, a_k \in \mathcal{S}_{N_i}(\mathbb{C})^+, \epsilon_k \in \mathbb{R}$ s.t.
 $\forall z \in \mathbb{C}^+,$

$$\Sigma_i(z) = C_i + \sum_{k=1}^{L_i} \frac{1}{z - \epsilon_k} a_k \quad (10)$$

Then $\forall i = 1, N, \Delta_i$ is well-defined and there exists $\tilde{L}_i \in \mathbb{N},$
 $\forall k = 1, \tilde{L}_i, \tilde{a}_k \in \mathcal{S}_{N_i}(\mathbb{C})^+, \tilde{\epsilon}_k \in \mathbb{R}$ s.t. $\forall z \in \mathbb{C}^+,$

$$\Delta_i(z) = \sum_{k=1}^{\tilde{L}_i} \frac{1}{z - \tilde{\epsilon}_k} \tilde{a}_k \quad (11)$$

Issue : $\tilde{L}_i > L_i$, no finite bath solution with IPT

Self-coherent equation well-posedness for any bath

Extension : $\Sigma(z) = C + \int_{\mathbb{R}} \frac{1}{z-\epsilon} d\mu(\epsilon)$, μ a $\mathcal{S}(\mathbb{C})^+$ -valued measure

Proposition : Self-coherent infinite bath well-posedness

Assume $\forall i = 1, N$, $\exists C_i \in \mathcal{S}_{N_i}(\mathbb{C})^+$ and μ_i a $\mathcal{S}_{N_i}(\mathbb{C})^+$ -valued measure (with integrability condition), s.t. $\forall z \in \mathbb{R}$,

$$\Sigma_i(z) = C_i + \int_{\mathbb{R}} \frac{1}{z-\epsilon} d\mu_i(\epsilon) \quad (12)$$

Then $\forall i = 1, N$, Δ_i is well-defined and there exists ν_i a *finite* $\mathcal{S}_{N_i}(\mathbb{C})^+$ -valued measure such that

$$\Delta_i(z) = \int_{\mathbb{R}} \frac{1}{z-\epsilon} d\nu_i(\epsilon) \quad (13)$$

In the literature : only found with $N = |\Lambda|$ (one site per impurity), with Γ the Gibbs state at β, μ , using *Matsubara's* Green's functions and frequencies $\omega_n = \frac{\pi(2n+1)}{\beta}$
Given Δ , IPT proceeds in two steps :

In the literature : only found with $N = |\Lambda|$ (one site per impurity), with Γ the Gibbs state at β, μ , using *Matsubara's* Green's functions and frequencies $\omega_n = \frac{\pi(2n+1)}{\beta}$
Given Δ , IPT proceeds in two steps :

- First compute $\forall n \in \mathbb{N}$,

$$\Sigma_n = \frac{U}{2} + U^2 \int_0^\beta e^{i\omega_n \tau} \left(\frac{1}{\beta} \sum_{n' \in \mathbb{Z}} e^{-i\omega_{n'} \tau} \frac{1}{i\omega_{n'} - h_{\text{imp}} + \mu - \Delta(i\omega_{n'})} \right)^3 d\tau \quad (14)$$

In the literature : only found with $N = |\Lambda|$ (one site per impurity), with Γ the Gibbs state at β, μ , using *Matsubara's* Green's functions and frequencies $\omega_n = \frac{\pi(2n+1)}{\beta}$
Given Δ , IPT proceeds in two steps :

- First compute $\forall n \in \mathbb{N}$,

$$\Sigma_n = \frac{U}{2} + U^2 \int_0^\beta e^{i\omega_n \tau} \left(\frac{1}{\beta} \sum_{n' \in \mathbb{Z}} e^{-i\omega_{n'} \tau} \frac{1}{i\omega_{n'} - h_{\text{imp}} + \mu - \Delta(i\omega_{n'})} \right)^3 d\tau \quad (14)$$

- "Analytic continue" it : find $\Sigma : \mathbb{C}_+ \rightarrow \mathbb{C}_+$ analytic s.t.

$$\forall n \in \mathbb{N}, \Sigma(i\omega_n) = \Sigma_n \quad (15)$$

Proposition : IPT first-step well-posedness

Given ν a positive measure (w. integrability conditions) s.t. $\forall z \in \mathbb{C}_+$,

$$\Delta(z) = \int_{\mathbb{R}} \frac{1}{z - \epsilon} d\nu(\epsilon) \quad (16)$$

Then $\forall n \in \mathbb{N}$, Σ_n is well-defined and $\exists \mu$ a positive and *finite* measure s.t. $\forall n \in \mathbb{N}$,

$$\Sigma_n = \frac{U}{2} + \int_{\mathbb{R}} \frac{1}{i\omega_n - \epsilon} d\mu(\epsilon) \quad (17)$$

Proposition : IPT second-step well-posedness

The analytical continuation problem

$$\text{Find } \Sigma : \mathbb{C}_+ \rightarrow \mathbb{C}_+ \text{ analytic s.t. } \forall n \in \mathbb{N}, \Sigma(i\omega_n) = \Sigma_n \quad (18)$$

admits a unique solution if

- Δ represents a bath of finite dimension
- ν is compactly supported

Solution : ν finite, (probably not compactly supported) ...

Ongoing :

- Necessary conditions on the solution.
- Functional equation on the *density* of ν if $\nu \ll \text{Lebesgue}$.
- Convergence study : *which topology* ?

Ongoing :

- Necessary conditions on the solution.
- Functional equation on the *density* of ν if $\nu \ll \text{Lebesgue}$.
- Convergence study : *which topology* ?



(a) S. Perrin-Roussel



(b) É. Cancès



(c) M. Vinteler
(M1 student)