Nuclear superfluidity The Pairing Hamiltonian as a many-body testbed

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Workshop on model systems in quantum mechanics January 12th, 2024



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Outline



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Nuclear phenomenology

- Odd-even staggering of experimental binding energies along isotopic chains
- Three-point mass differences give estimate for the (neutron) pairing gap

$$\Delta_N^{(3)} = \frac{(-1)^N}{2} \left(E_{N+1} - 2E_N + E_{N-1} \right)$$

• Experimental evidence of formation of Cooper pairs in atomic nuclei

short-range attractive two-body interaction

• Nuclear phenomenology emerges from interplay of pairing and deformation

→ see also talk by **D.** Lacroix



Nuclear masses of tin isotopes (AME2020)

• One-parameter interaction describing superfluidity

$$H_{\text{pairing}} = \sum_{p} \epsilon_{p} \left(c_{p}^{\dagger} c_{p} + c_{\bar{p}}^{\dagger} c_{\bar{p}} \right) + g \sum_{pq} c_{p}^{\dagger} c_{\bar{p}}^{\dagger} c_{\bar{q}} c_{q}$$

• Generation of a pair of time-reversed states

$$|p\rangle = |n_p l_p j_p m_p\rangle \quad \rightarrow \quad |\bar{p}\rangle = |n_p l_p j_p - m_p\rangle$$





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 Admits for a transition to superfluid regime at critical coupling strength

Number of levels
$$\Omega = 8$$
Occupied levels $N_{\rm OCC} = 4$



One-parameter interaction describing superfluidity



• Richardson solution: exact wave function is written from pair creation operators

Richardson, PL (1965), PR (1966)

$$|\Psi\rangle = B_1^{\dagger} \cdot \ldots \cdot B_{\Omega}^{\dagger}|0\rangle \qquad \qquad B_{\alpha}^{\dagger} = \sum_{p=1}^{\Omega} \frac{1}{2\epsilon_p - E_{\alpha}} P_p^{\dagger}$$

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• Solving coupled system of equations provides unknown pair energies E_a

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- Modifications are required in case of unpaired particles (non-zero seniority)
- Extend full configuration interaction (FCI) capacities: limited to ~20 levels

dynamic correlations

(expansion on top of dominant Slater determinant)



static correlations (collectivity through symmetry projection)

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• BCS wave-function ansatz for superfluid system

(Bardeen-Cooper-Schrieffer)

$$|\Phi_{\text{BCS}}\rangle = \prod_{p>0} \left(u_p + v_p c_p^{\dagger} c_{\bar{p}}^{\dagger} \right) |0\rangle$$

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Solution obtained from solving BCS equations

$$v_{\rho}^{2} = \frac{1}{2} \left(1 - \frac{\xi_{\rho}}{\sqrt{\xi_{\rho}^{2} + \Delta^{2}}} \right)$$
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$$\Delta = \langle \Phi_{\text{BCS}} | c_{\rho}^{\dagger} c_{\bar{\rho}}^{\dagger} | \Phi_{\text{BCS}} \rangle = g \sum_{\rho > 0} u_{\rho} v_{\rho}$$

BCS gap parameter

• HF theory recovered in the limit of vanishing gap



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- Application of particle-number projector restores the broken symmetry

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• Similar expressions hold for other symmetries, e.g., rotational invariance SU(2)

Lacroix, Gambacurta, PRC (2012)

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• Textbook formulation: Hartree-Fock Slater determinant as reference state

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i: holes (occupied) *a*: particles (unoccupied)

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Quasiparticle formulation: MBPT expansion around symmetry-broken BCS vacuum

 $E^{(2)} = -\frac{1}{2} \sum_{pq} g^2 \frac{\left(u_p^2 v_q^2 + u_q^2 v_p^2\right)^2}{E_p + E_q} \qquad E_p > 0: \text{ quasi-particle energies}$

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• **Projected MBPT:** expectation-value formulation with particle-number projector

$$E = \frac{\langle \Psi | P_A^{\dagger} H P_A | \Psi \rangle}{\langle \Psi | P_A | \Psi \rangle} \qquad \qquad |\Psi\rangle = |\Phi_{BCS}\rangle + |\Psi^{(1)}\rangle$$

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• Standard HF-MBPT is recovered in limiting case of vanishing pairing gap

Results of perturbation theory



Henderson et al., PRC (2014)

• Bogoliubov coupled-cluster (BCC) theory: CC extension for BCS reference states

$$|\Psi_{\rm BCC}\rangle = e^{T}|\Phi_{\rm BCS}\rangle$$

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$$\beta_{\rho}^{\dagger} = u_{\rho}c_{\rho}^{\dagger} + v_{\rho}c_{\bar{\rho}} \qquad \qquad \beta_{\bar{\rho}}^{\dagger} = u_{\rho}c_{\bar{\rho}}^{\dagger} - v_{\rho}c_{\rho}$$

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$$T_2 = \frac{1}{2} \sum_{pq} t_{pq} \mathcal{P}_p^{\dagger} \mathcal{P}_q^{\dagger} \qquad \qquad \mathcal{P}_p^{\dagger} = \beta_p^{\dagger} \beta_{\bar{p}}^{\dagger}$$

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$$0 = \langle \Phi_{\text{BCS}} | \mathcal{P}_p \mathcal{P}_q e^{-T} (H - \lambda A) e^T | \Phi_{\text{BCS}} \rangle \quad \forall p, q$$

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• Projected coupled-cluster energy in presence of projection operator

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Qiu et al., PRC (2019) Duguet, JPG (2014)

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- Symmetry projection beyond mean-field constitutes a highly non-trivial task
- Numerical implementation is computationally very demanding in practice





Large-scale ab initio applications



Particle-number-broken many-body frameworks for open-shell nuclei

No full projection yet!

• Question: can we emulate the many-body solution based on training data?



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training data	training vectors		target	coupling
$\{g_1,, g_N\}$	$\{ \Psi(g_1)\rangle, \dots, \Psi(g_N)\}$	>	g_{\odot}	 Ψ(g_☉) ⟩

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Hot topic: emerging field in nuclear physics with numerous applications

Duguet et al., arXiv:2310.19419 (2023)

Performance of the EC emulator



Companys Franzke, Tichai, Hebeler, Schwenk, PRC (2024)

- Introduction of different training manifolds: normal, superfluid, mixed
- One-sided training manifolds are unreliable for non-trained regime
- Mixed training manifold yields consistent prediction for all couplings

Selection of appropriate training vectors crucial!

similar study using DMRG:

Baran, Nichita, PRB (2023)

Performance of the EC emulator



Conclusions

Pairing Hamiltonian as many-body testbed

- Integrable Richardson solution for arbitrary system size
- Emergence of critical coupling separating normal and superfluid phase
- Breakdown of conventional many-body expansions

Lesson: Hartree-Fock-based schemes are doomed to fail

Implications on many-body frameworks

- Capture static correlations via spontaneous symmetry breaking
- Account for dynamic correlations using many-body expansion
- Parametric dependence can be efficiently emulated

Lesson: reference state must capture important static correlations

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Lesson: Hartree

Take-home message: Schematic models reveal key correlations relevant for realistic applications.

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