

Harmonium: a model to study approximations in TDDFT

M.-S. Dupuy¹

¹Laboratoire Jacques-Louis LIONS, Sorbonne-Université, Paris

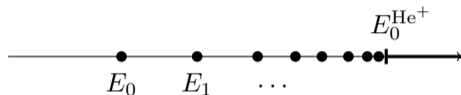
Model systems in quantum mechanics 2024

Nonrelativistic Helium atom

$$\mathcal{H}_{\text{He}} = -\frac{1}{2}\Delta_{r_1} - \frac{1}{2}\Delta_{r_2} - \frac{Z}{|r_1|} - \frac{Z}{|r_2|} + \frac{1}{|r_1 - r_2|},$$

acting on $L^2(\mathbb{R}^3 \times \{\uparrow, \downarrow\}) \wedge L^2(\mathbb{R}^3 \times \{\uparrow, \downarrow\})$ with domain $H^2((\mathbb{R}^3 \times \{\uparrow, \downarrow\})^2)$ ($Z = 2$).

Drawback ground-state not explicitly known



Spectrum of \mathcal{H}_{He}

Harmonium or Hooke's atom

$$\mathcal{H}_{\text{harm}} = -\frac{1}{2}\Delta_{r_1} - \frac{1}{2}\Delta_{r_2} + \frac{1}{2}k|r_1|^2 + \frac{1}{2}k|r_2|^2 + \frac{1}{|r_1 - r_2|},$$

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Properties

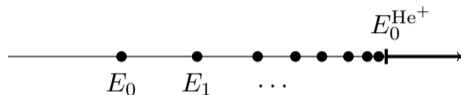
- ▶ for some values of k , the ground-state is explicit
- ▶ $\mathcal{H}_{\text{harm}}$ has only a discrete spectrum with eigenvalues $E_k \rightarrow \infty$ (potential tends to ∞)

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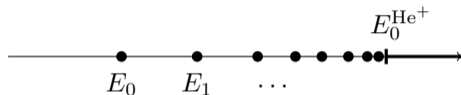
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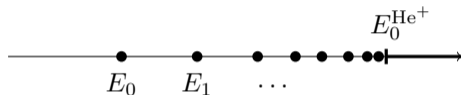
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Calculation of the ground-state

Change of variables $R = \frac{1}{2}(r_1 + r_2)$ and $u = r_2 - r_1$

▶ separation of variables: $\mathcal{H}_{\text{harm}} = \mathcal{H}_c(u) + \mathcal{H}_{\text{ho}}(R)$ with $\begin{cases} \mathcal{H}_{\text{ho}}(R) = -\frac{1}{4}\Delta_R + k|R|^2, \\ \mathcal{H}_c(u) = -\Delta + \frac{k}{4}|u|^2 + \frac{1}{|u|}. \end{cases}$

▶ eigenfunctions: $\Psi(r_1, r_2) = \chi\left(\frac{r_1+r_2}{2}\right)\Phi(r_2 - r_1)$ with $\mathcal{H}_{\text{ho}}\chi = E_{\text{ho}}\chi$ and $\mathcal{H}_c\Phi = E_c\Phi$

▶ eigenvalues $E = E_{\text{ho}} + E_c$.

Calculation of Φ

▶ $\frac{k}{4}|u|^2 + \frac{1}{|u|}$ is a central potential so $\Phi(u) = R_\ell(|u|)Y_{\ell m}(\hat{u})$

▶ look for $R_\ell(u) = \sum_{j=0}^{\infty} a_j u^j \exp(-\frac{\sqrt{k}}{4}u^2)$

▶ for $k = \frac{1}{4}$, $\sum_{j=0}^{\infty} a_j u^j = 1 + \frac{u}{2}$

Ground-state of the Harmonium ($k = \frac{1}{4}$)

$$\Psi(r_1, r_2) = \frac{1}{2\sqrt{8\pi^{5/2} + 5\pi^3}} \left(1 + \frac{1}{2}|r_1 - r_2|\right) \exp\left(-\frac{|r_1|^2 + |r_2|^2}{4}\right).$$

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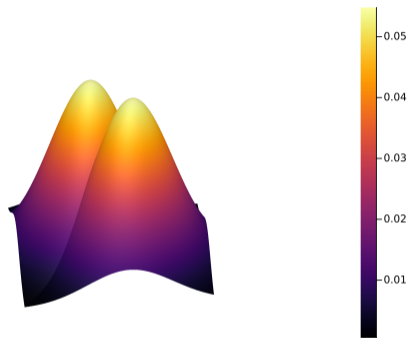
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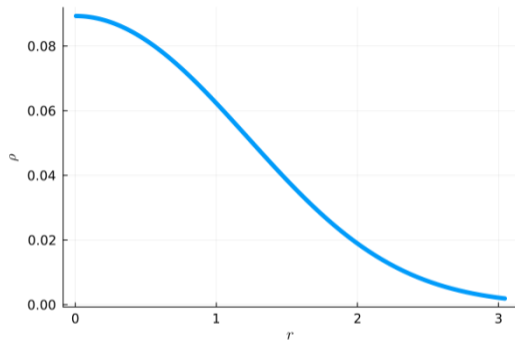
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Plots of the ground-state and the electronic density



(a) Wave function $\Psi(r_1, r_2)$



(b) Electronic density ρ

Remarks

- ▶ electron-electron cusp in the Harmonium wave function
- ▶ no nucleus-electron cusp in the wave function or the electronic density

LR TDDFT “theorem”

There is an xc-kernel f_{xc} such that the many-body one-body linear response function χ is given by

$$\widehat{\chi}(\omega) = \widehat{\chi}_0(\omega) + \widehat{\chi}_0(\omega) \widehat{f}_{xc}(\omega) \widehat{\chi}(\omega),$$

where χ_0 is the linear response function of the noninteracting Kohn-Sham Hamiltonian.

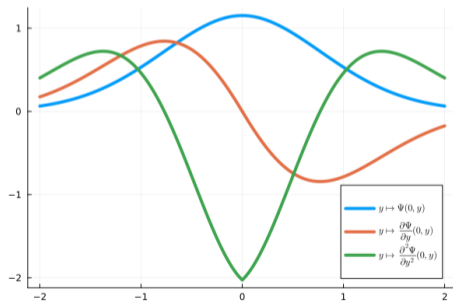
Tests on a one-dimensional Harmonium

$$\mathcal{H} = -\frac{1}{2} \frac{\partial^2}{\partial x^2} - \frac{1}{2} \frac{\partial^2}{\partial y^2} + \frac{1}{2} k(x^2 + y^2) - |x - y|,$$

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Properties

- ▶ spectrum of \mathcal{H} is only discrete
- ▶ all the eigenfunctions are explicit (via the Whittaker function U)
- ▶ electron-electron cusp in the 2nd derivative



Ground-state of \mathcal{H} and its derivatives

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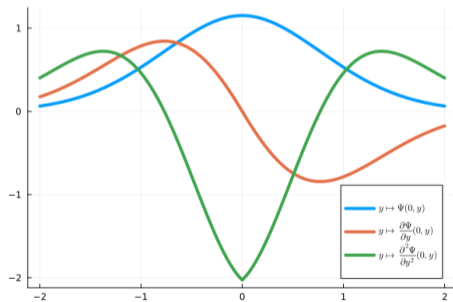
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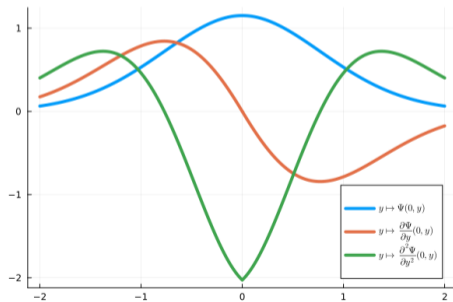
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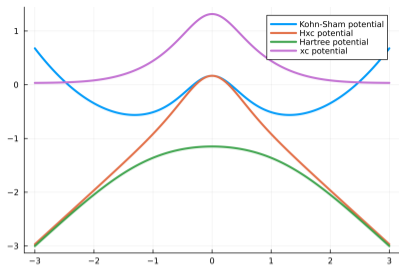
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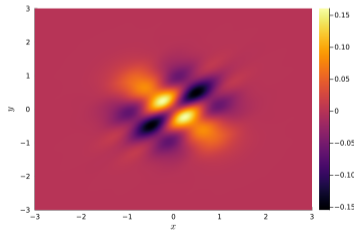
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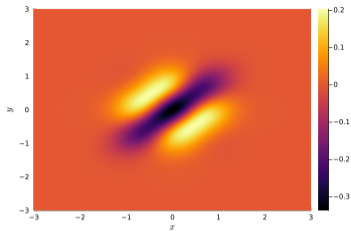
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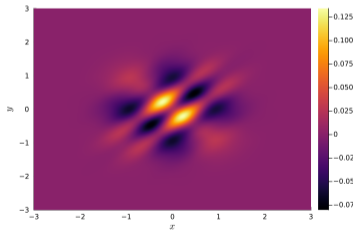
Kohn-Sham and other potentials



Difference $\hat{\chi}(x, y; 0) - \hat{\chi}_0(x, y; 0)$



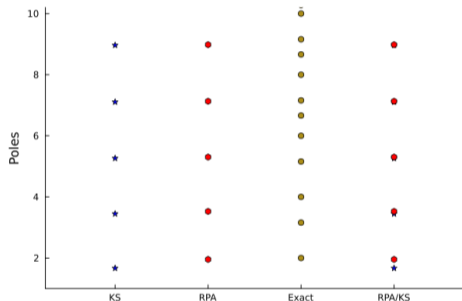
Exact linear response function $\hat{\chi}(x, y; 0)$



Difference $\hat{\chi}(x, y; 0) - \hat{\chi}^{\text{RPA}}(x, y; 0)$

Poles of the linear response function $\hat{\chi} \equiv \text{excited energies } E_k - E_0 \ (k \geq 1)$

- ▶ TDDFT Dyson equation: correction to the KS eigenvalues



Poles of the Kohn-Sham LRF $\hat{\chi}_0$, the RPA LRF $\hat{\chi}^{\text{RPA}}$ and the exact LRF $\hat{\chi}$

Observations

- ▶ many more poles in the exact than RPA/KS linear response functions (failure of adiabatic xc kernels)
- ▶ RPA poles smaller than KS poles (positivity of the xc kernel)

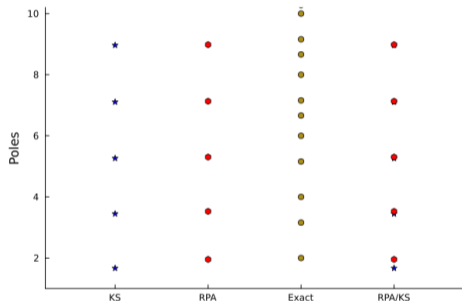
Reference for RPA pole shifting

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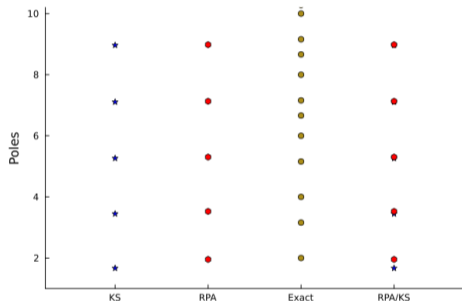
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